

Thermo-acoustic Oscillation of in a Closed Tube by Numerical Simulation

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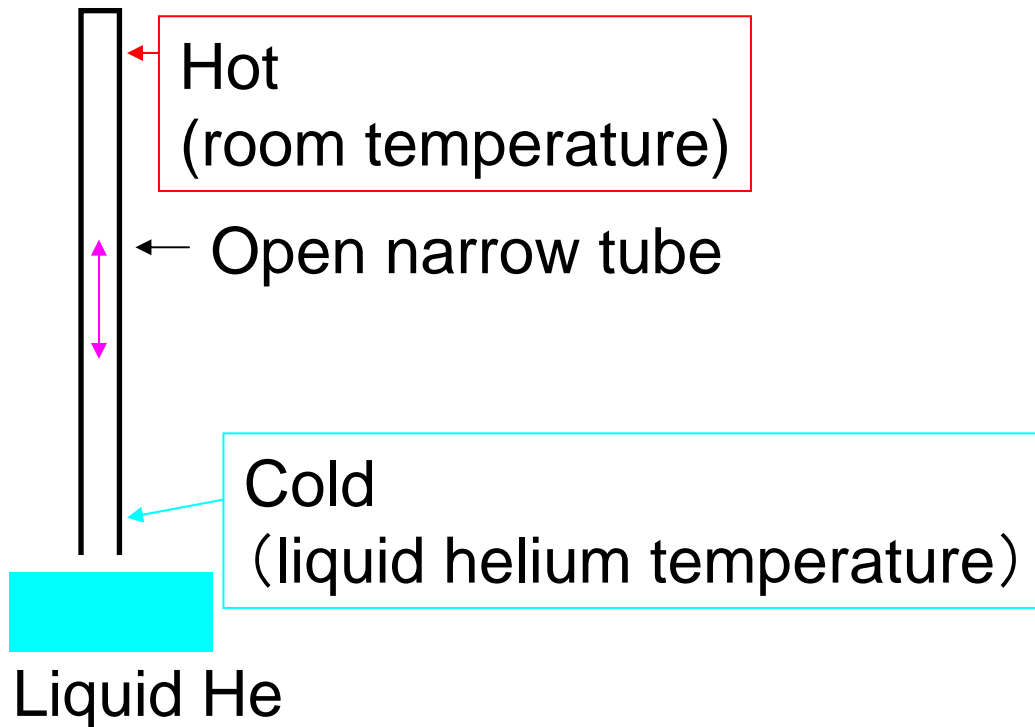
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1.1 Taconis oscillation

Taconis oscillation is one of the thermo-acoustic oscillations.



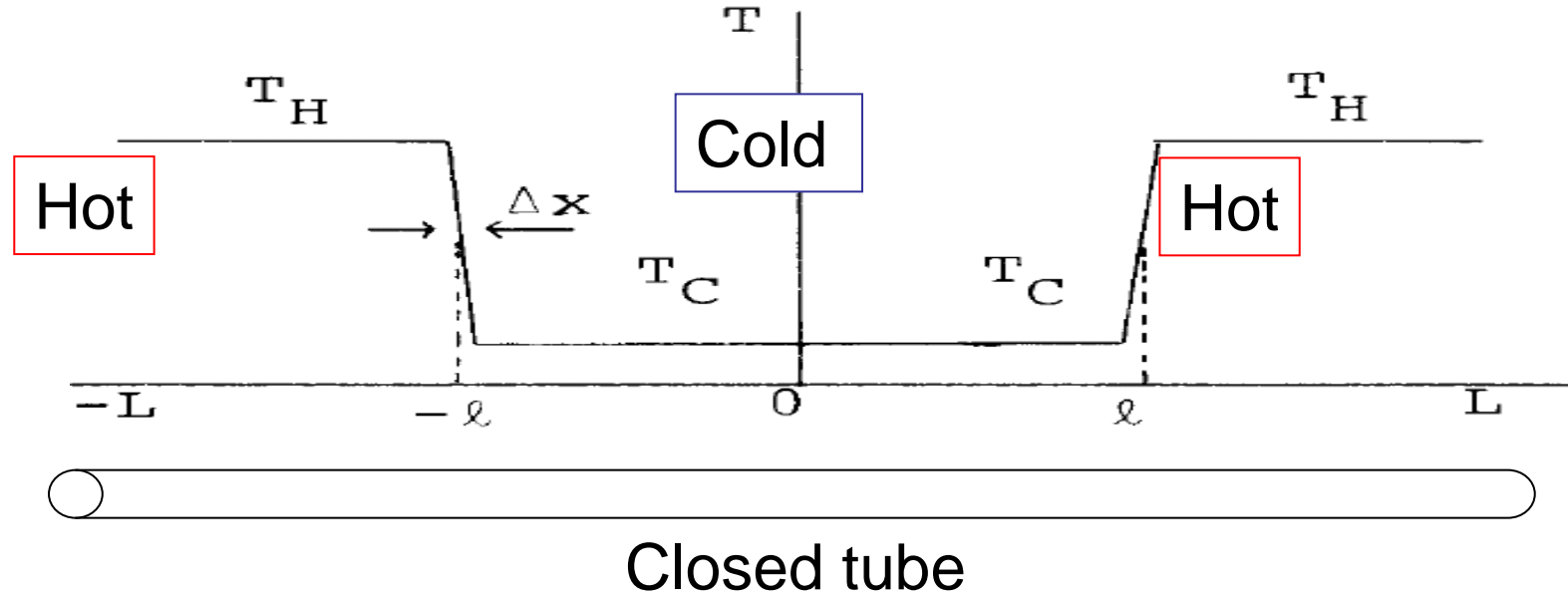
Experimental study
- Taconis et al (1949)

Theoretical study
- Kramers (1949)
- Rott (1969, 1973)
- Sugimoto et al (2006-2008)

The tube has large temperature gradient axially.
The spontaneous oscillation is observed in the tube.

1.2 Experiment by Yazaki et al

(J.Low.Temp.Phys., **41**, 1980)



Yazaki et al observed a standing wave for different large temperature ratios T_H/T_C .

They explained these phenomena by using the theory of Taconis oscillation in a open-closed tube.

1.3 Thermo-acoustic Engine

G.B. Chen, T. Jin (Cryogenics, **39**, 1999, 843-846)

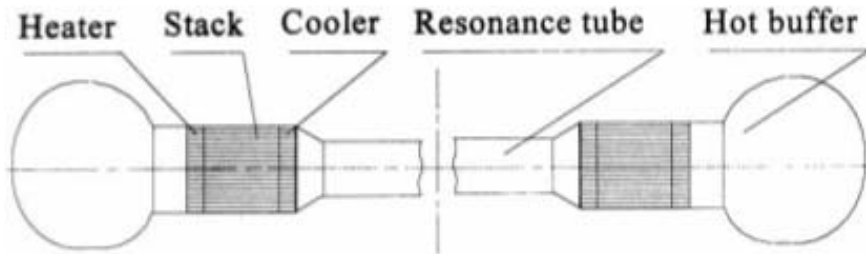
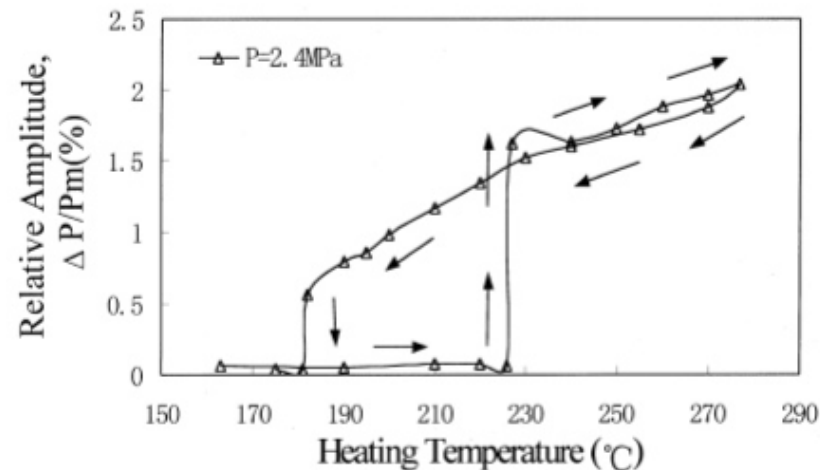


Fig. 1. Schematic of thermoacoustic engine.

The Pressure amplitude is measured when the heating temperature is changed.



The hysteresis phenomenon is observed.

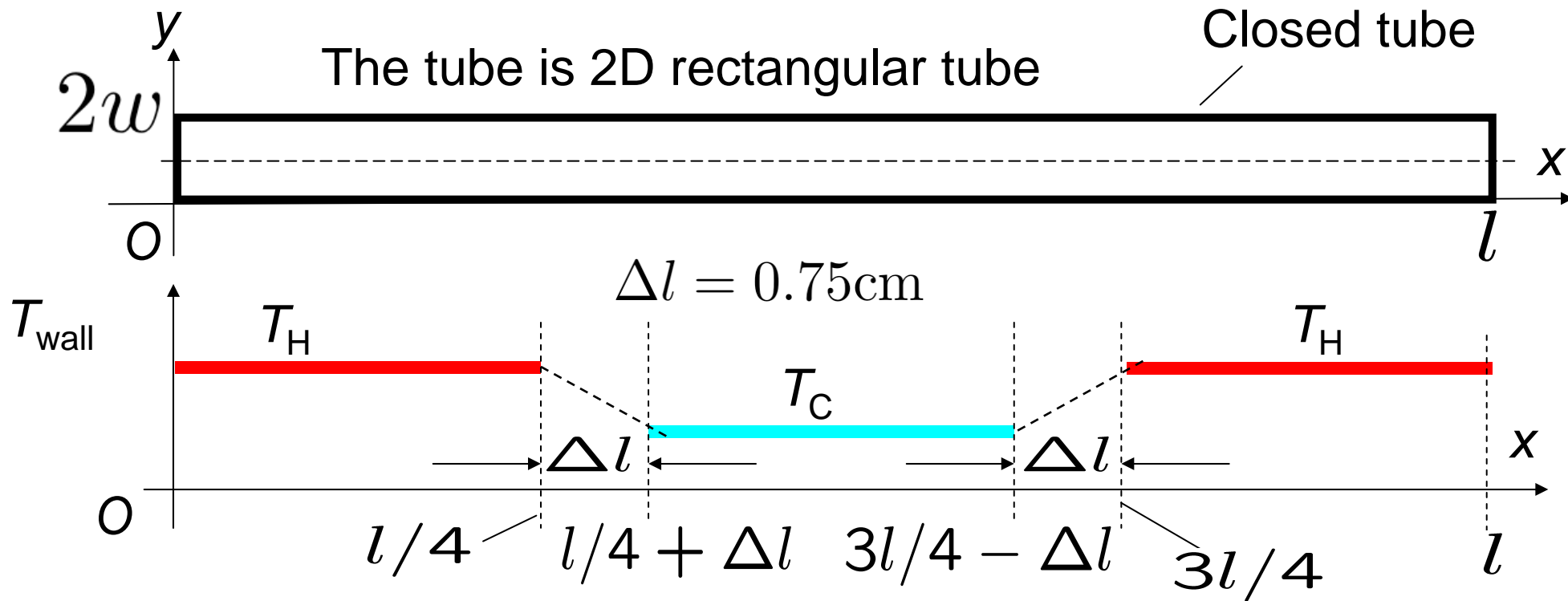


1.4 Objective

We observed **hysteresis phenomena in Taconis oscillation** which has not been reported experimentally.

The objectives are to analyze the Taconis oscillation in a closed tube using the numerical simulations, and to investigate the mechanism of hysteresis phenomenon.

2.1 Geometry



The fluid in the tube is **gaseous helium** at the room temperature $T_H=300$ K and the pressure $p_0 = 0.175 \times 10^5$ Pa, 1×10^5 Pa. (We use different initial pressures p_0 in order to obtain the different thickness of the boundary layers.)

2.2 Basic equations

2D compressible Navier-Stokes equations

$$\partial_t \mathbf{q} + \partial_x \mathbf{E} + \partial_y \mathbf{F} = 1/Re(\partial_x \mathbf{R} + \partial_y \mathbf{S})$$

$$\mathbf{q} = \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ e \end{pmatrix}$$

$$\mathbf{E} = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ (e + p)u \end{pmatrix}$$

$$\mathbf{F} = \begin{pmatrix} \rho v \\ \rho uv \\ \rho v^2 + p \\ (e + p)v \end{pmatrix}$$

$$\mathbf{R} = \begin{pmatrix} 0 \\ \tau_{xx} \\ \tau_{xy} \\ R_4 \end{pmatrix}, \mathbf{S} = \begin{pmatrix} 0 \\ \tau_{xy} \\ \tau_{yy} \\ S_4 \end{pmatrix}$$

$$\tau_{xx} = 4/3\mu u_x - 2/3\mu v_y$$

$$\tau_{xy} = \mu(u_y + v_x)$$

$$\tau_{yy} = 4/3\mu v_y - 2/3\mu u_x$$

$$R_4 = u\tau_{xx} + v\tau_{xy} + \alpha\partial_x a^2$$

$$S_4 = u\tau_{xy} + v\tau_{yy} + \alpha\partial_y a^2$$

$$a^2 = \gamma(\gamma - 1)[e/\rho - 1/2(u^2 + v^2)]$$

$$\alpha = \frac{k}{Pr(\gamma - 1)}$$

ρ : density

u, v : velocity

e : total energy density

p : pressure

a : acoustic velocity

μ : viscosity

k : thermal conductivity

Pr : Prandtl number

γ : specific heat ratio

Stokes hypothesis

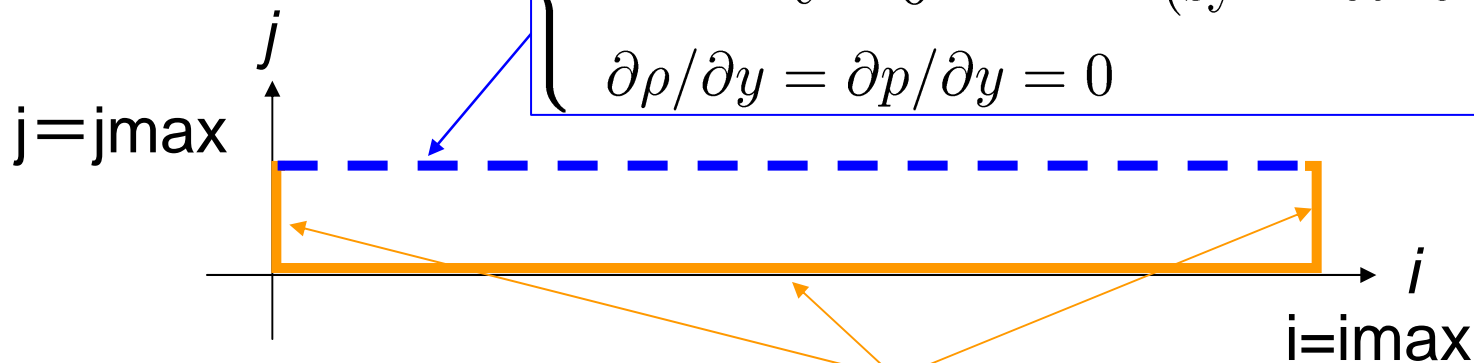
Sutherland law

2.3 Numerical calculation

- The block pentadiagonal matrix scheme
time development: 2nd-order accurate three-point backward scheme
convective terms: 4th -order accurate central differencing
viscous terms: 2nd -order accurate central differencing
- Grid 300 x 36

- Boundary conditions

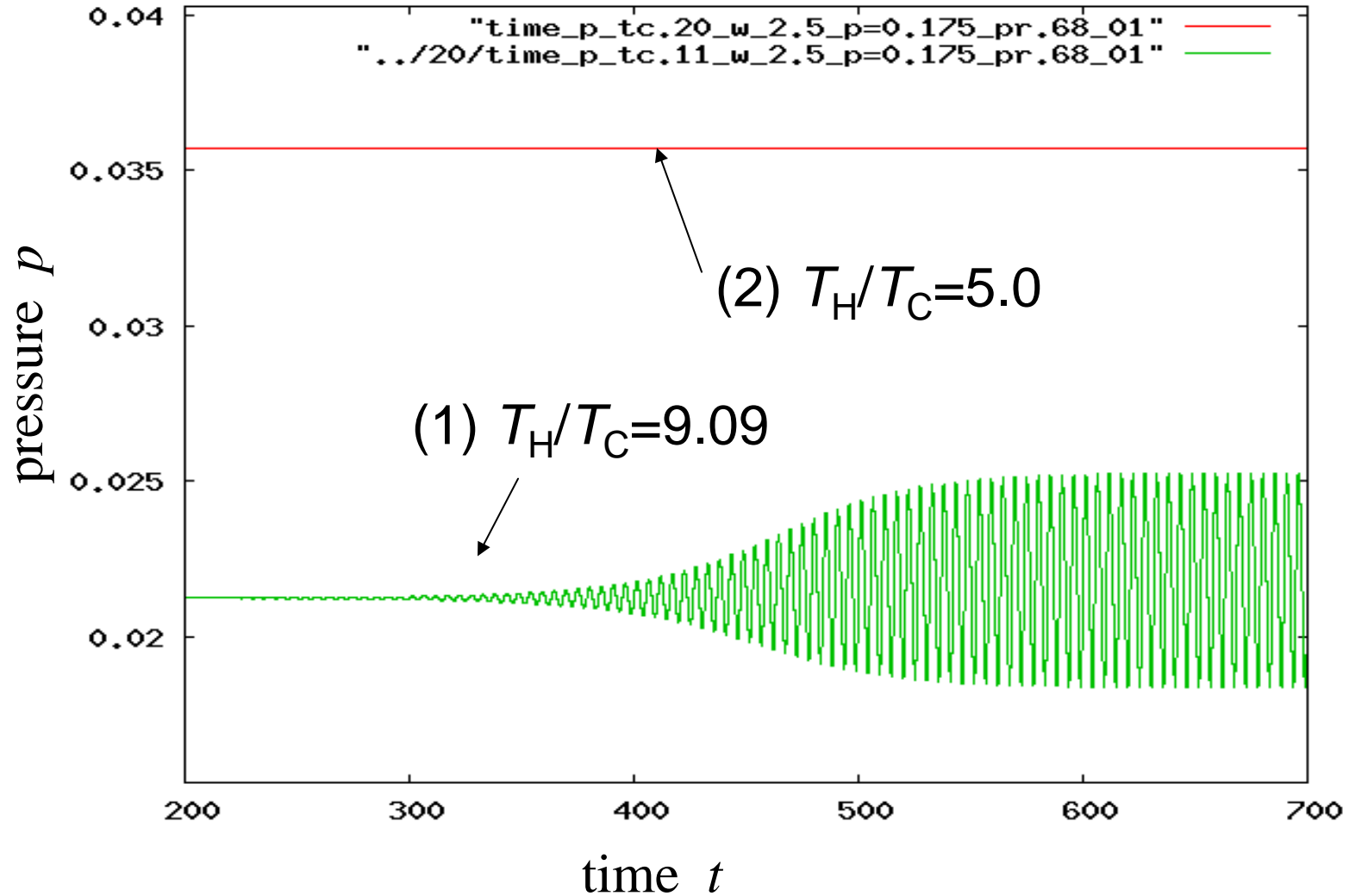
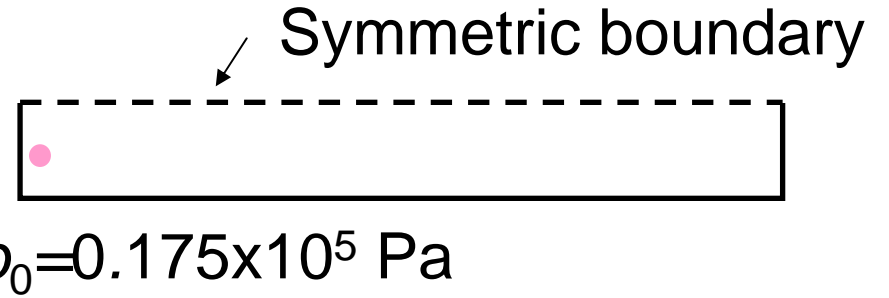
$$\left\{ \begin{array}{l} \partial u / \partial y = 0 \\ v = 0 \quad (\text{symmetric boundary}) \\ \partial \rho / \partial y = \partial p / \partial y = 0 \end{array} \right.$$



$$\left\{ \begin{array}{l} u = v = 0 \\ \partial p / \partial n = 0 \quad (\text{tube wall}) \\ T = T_{\text{wall}} \end{array} \right.$$

3.1 Time history of pressure

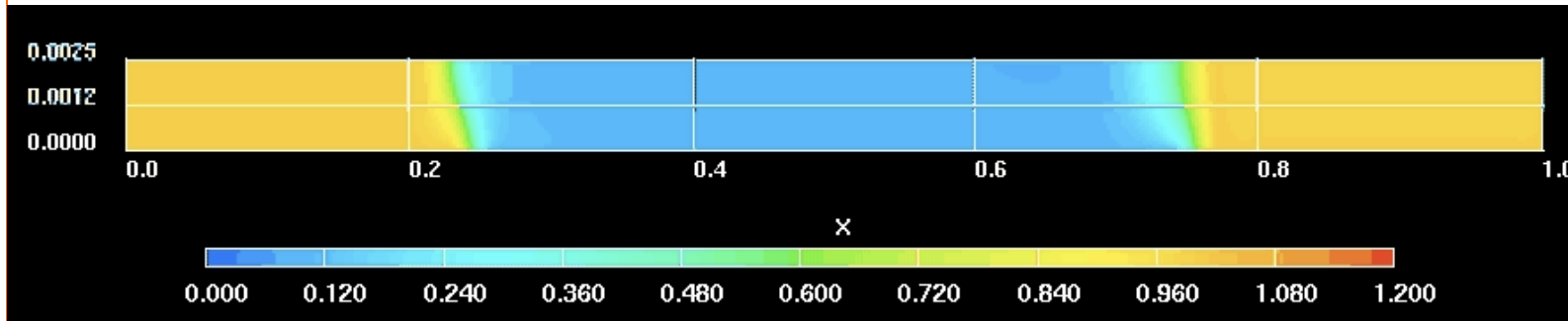
- (1) Steady oscillation state
- (2) Quiescent state



3.2 Thickness of thermal boundary layer

$$p_0 = 0.175 \times 10^5 \text{ Pa}$$

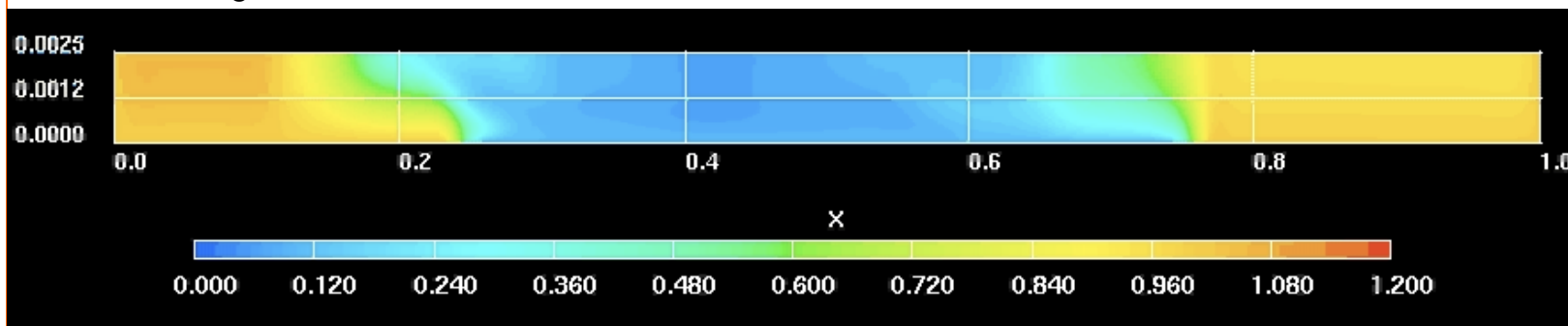
$$T_H/T_C = 9.09$$



$$\delta_{\alpha H}/w \approx 2.1$$

$$\delta_{\alpha C}/w \approx 0.4$$

$$p_0 = 1 \times 10^5 \text{ Pa}$$

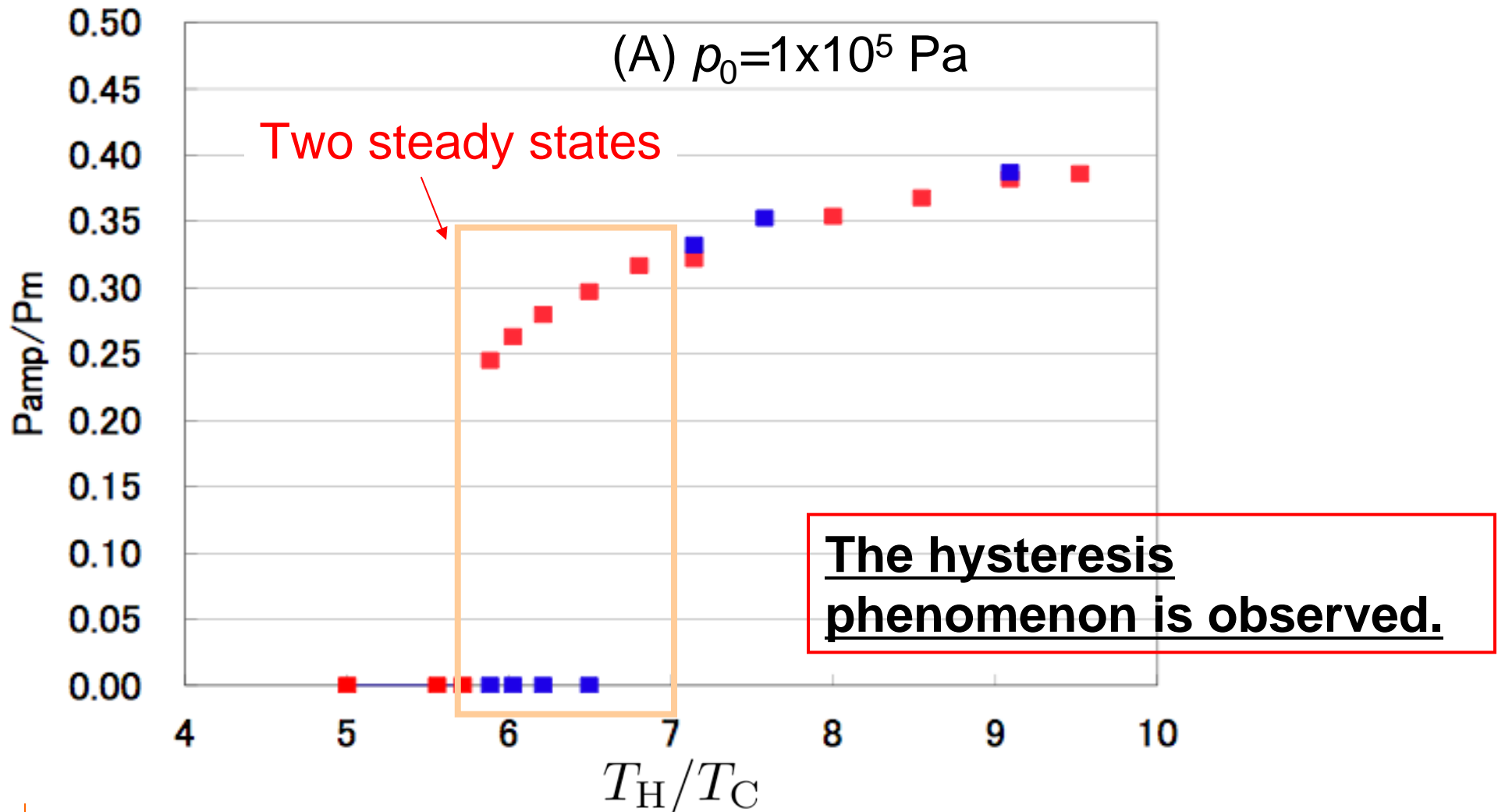


$$\delta_{\alpha H}/w \approx 1.0$$

$$\delta_{\alpha C}/w \approx 0.2$$

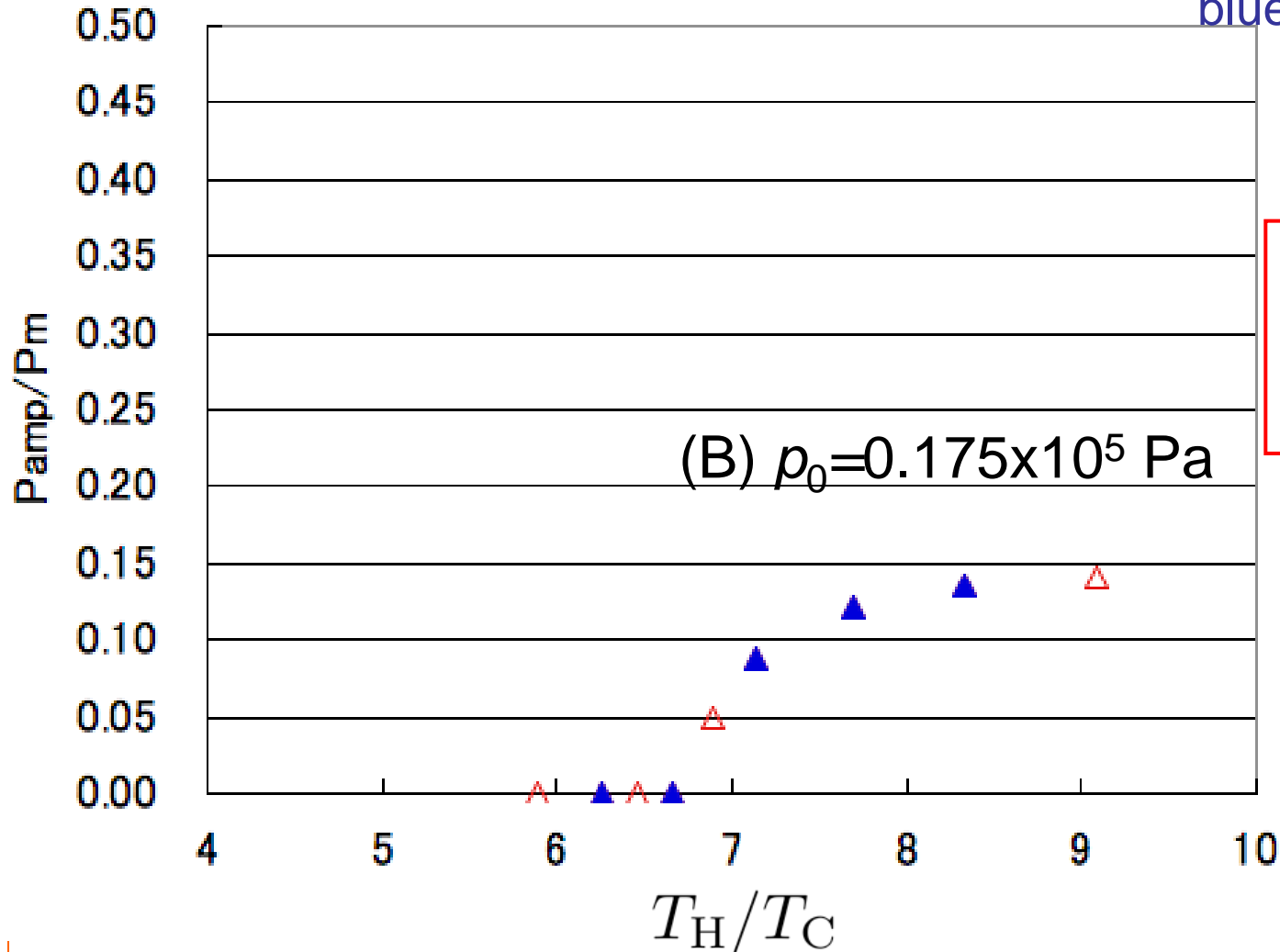
3.3 Pressure amplitude at $p_0=1 \times 10^5$ Pa

■: P_{amp} (A) $p_0=1 \times 10^5$ Pa
red: oscillation state initially
blue: quiescent state initially



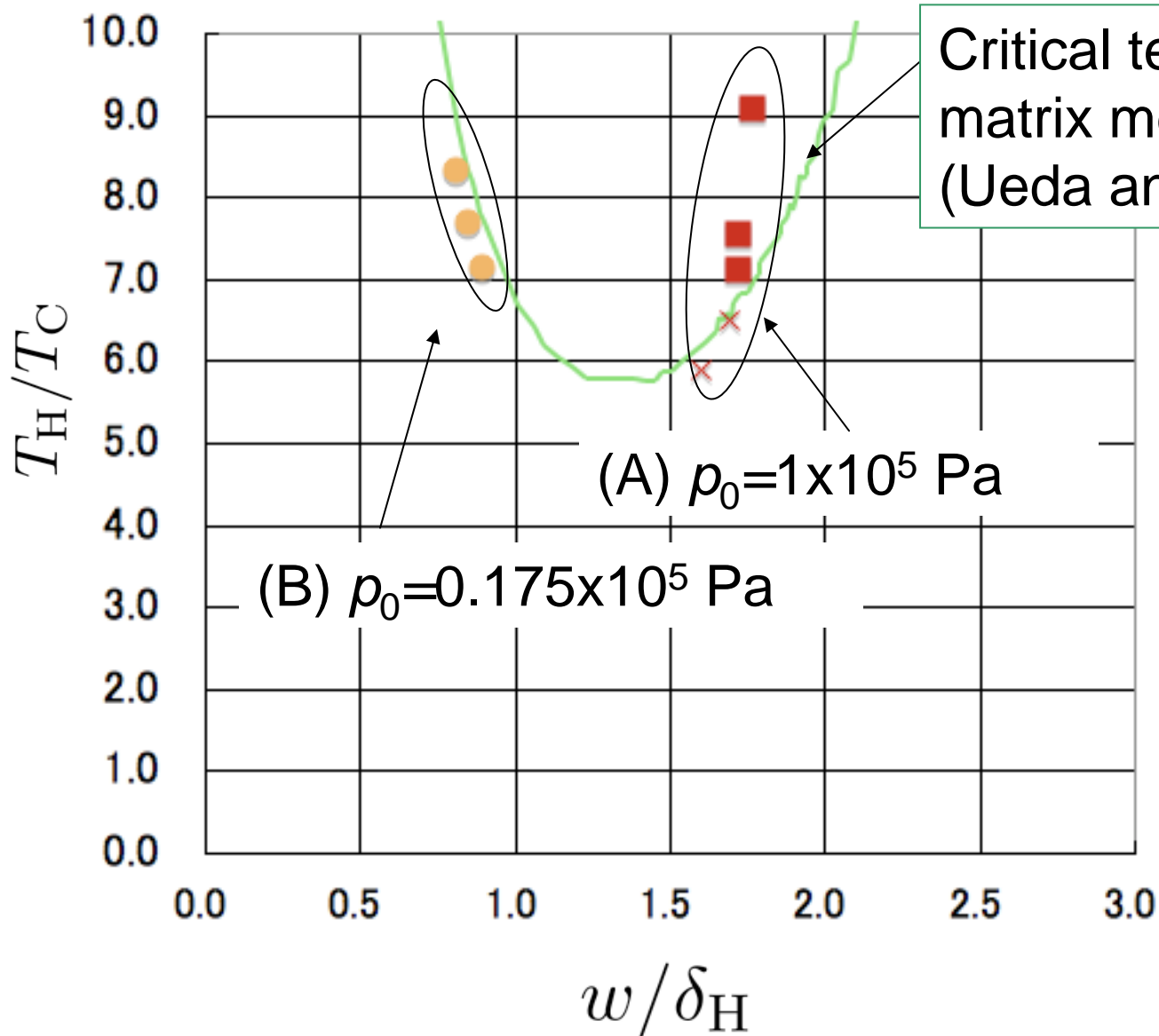
3.3 Pressure amplitude at $p_0=0.175 \times 10^5$ Pa

▲: P_{amp} (B) $p_0=0.175 \times 10^5$ Pa
red: oscillation state initially
blue: quiescent state initially



The hysteresis phenomenon is **not** observed.

3.4 Critical temperature ratio



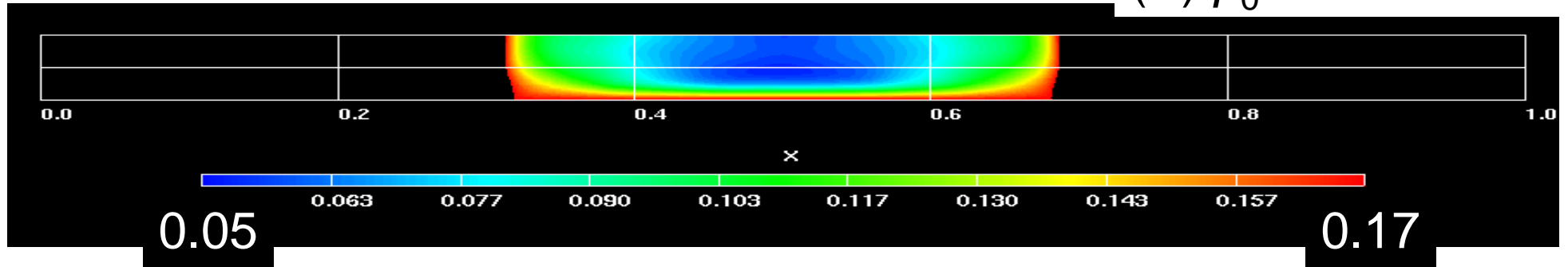
Critical temp. ratio by matrix method (Ueda and Kato 2008)

δ_H : thickness of the viscous boundary layer
 $\delta_H = (\nu_H/\Omega)^{0.5}$
 Ω : angular frequency
 ν_H : viscosity in hot part

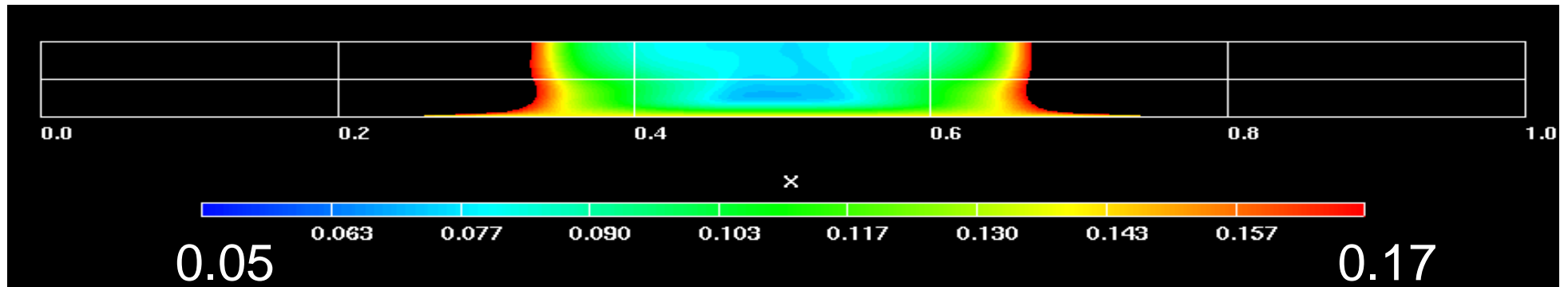
3.5 Time averaged temperature distribution

$T_H/T_C=5.88$ ($T_C=0.17$)

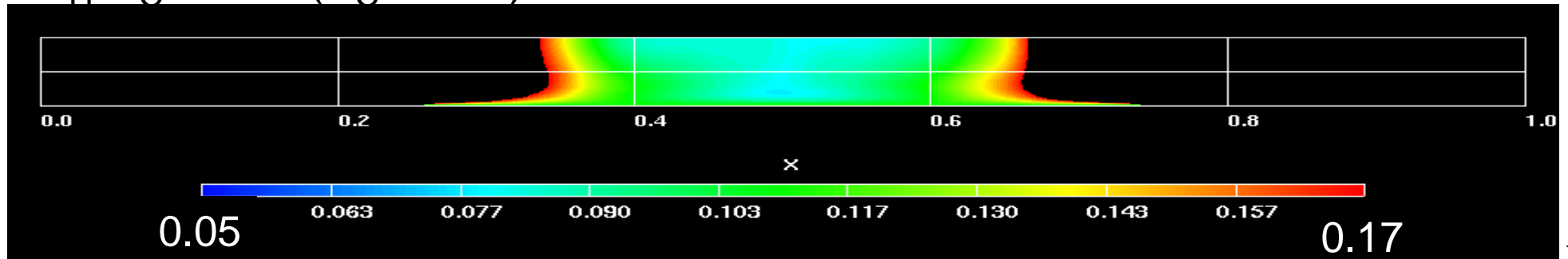
(A) $p_0=1 \times 10^5$ Pa



$T_H/T_C=7.14$ ($T_C=0.14$)



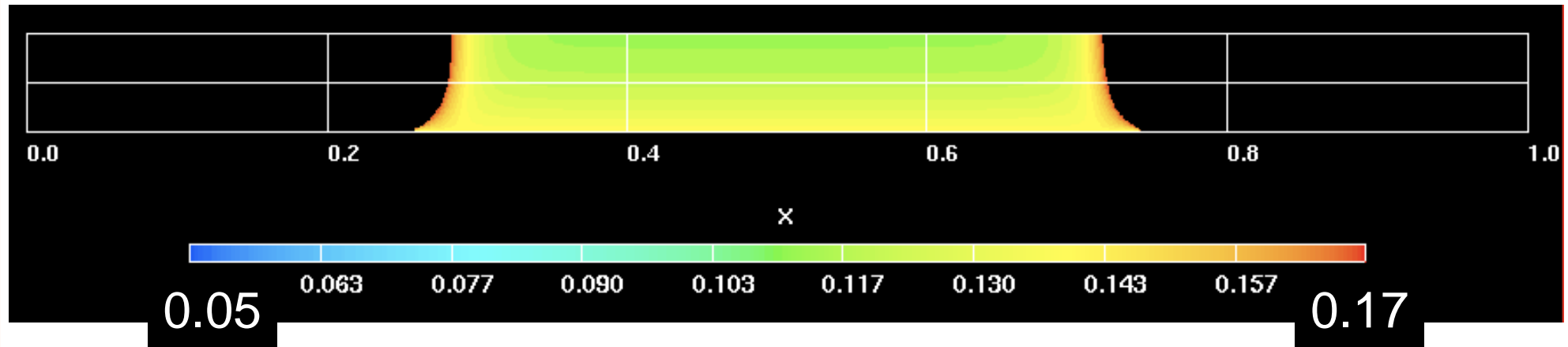
$T_H/T_C=9.09$ ($T_C=0.11$)



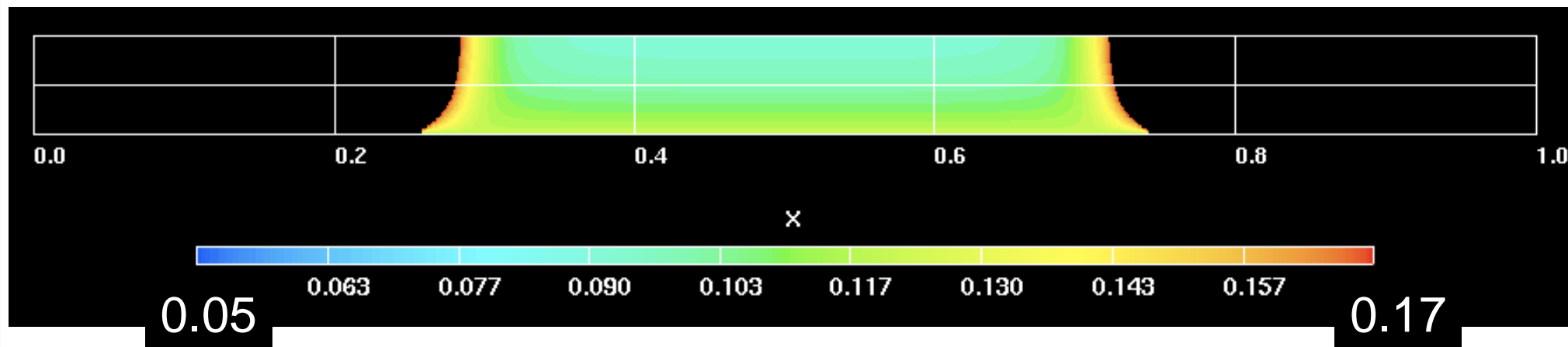
3.5 Time averaged temperature distribution

$$T_H/T_C=7.14 \quad (T_C=0.14)$$

$$(B) \quad p_0=0.175 \times 10^5 \text{ Pa}$$



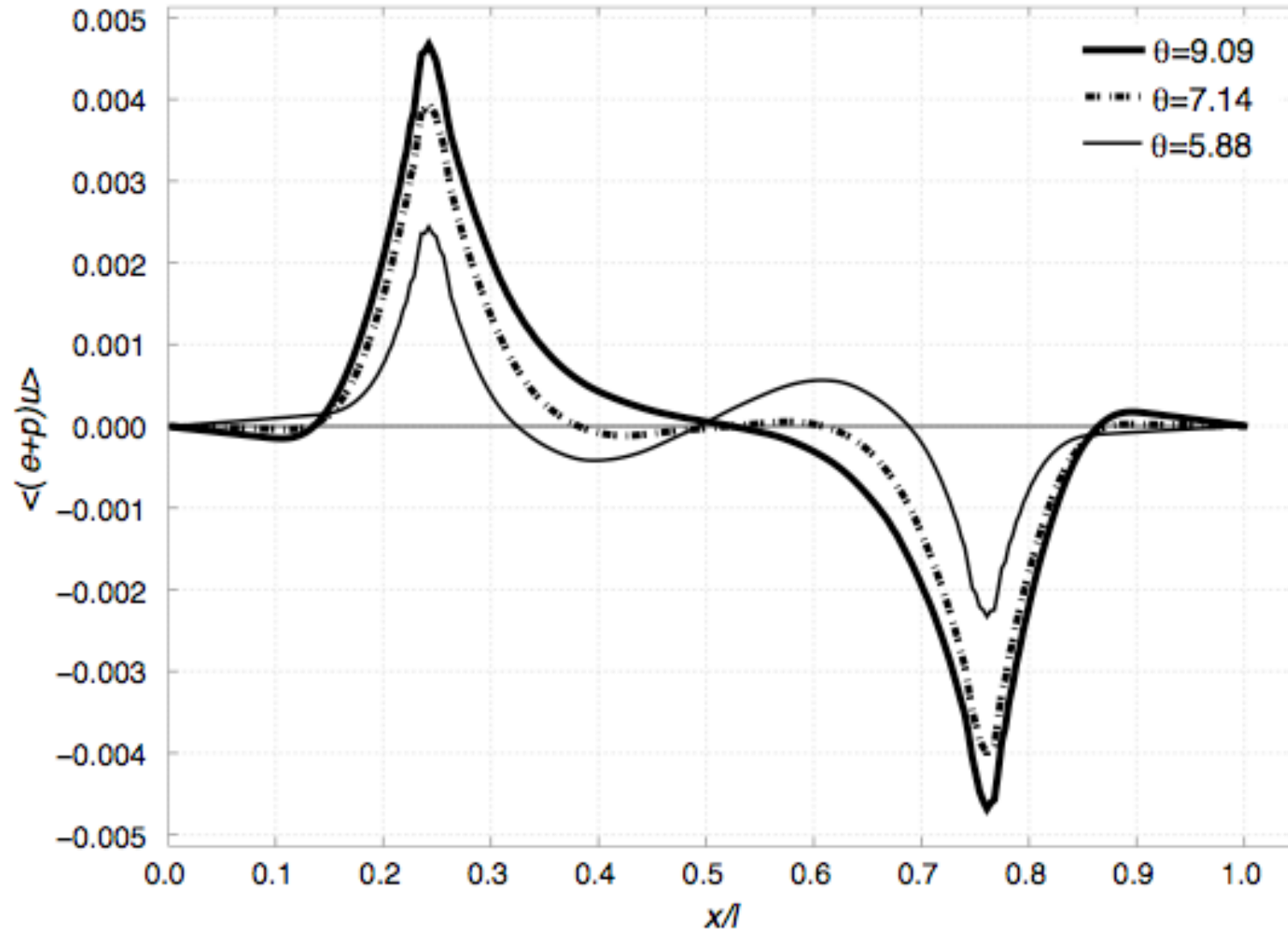
$$T_H/T_C=8.33 \quad (T_C=0.12)$$



3.6 Axial energy flux

(A) $p_0=1 \times 10^5$ Pa

$$\text{Axial energy flux} = (e + p)u$$



4 Summary

- We analyzed the thermo-acoustic oscillation (Taconis oscillation) in the 2D closed tubes by the numerical simulation.
- We observed hysteresis phenomena in the case of the initial pressure $p_0 = 1 \times 10^5$ Pa. On the other hand, we did not observe the hysteresis phenomenon in the case of $p_0 = 0.175 \times 10^5$ Pa.
- The heat pumping is observed in the cold region of the tube when the hysteresis phenomenon is observed. On the other hand, the effect of the heat pumping is hardly observed when the hysteresis phenomenon is not observed.